

# Sufficient Statistics

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April 11, 2019

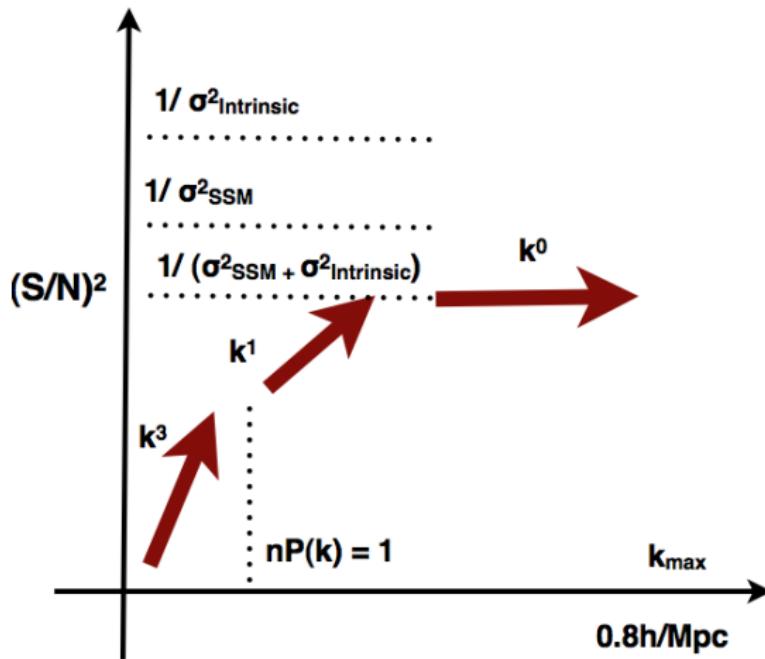
# Outline

- 1 Introduction to Sufficient Statistics
- 2 Prediction: PDF, the Log and  $A^*$  Power Spectrum, and Ising bias
- 3 Compactified Simulations

# Cosmological Information in LSS Surveys

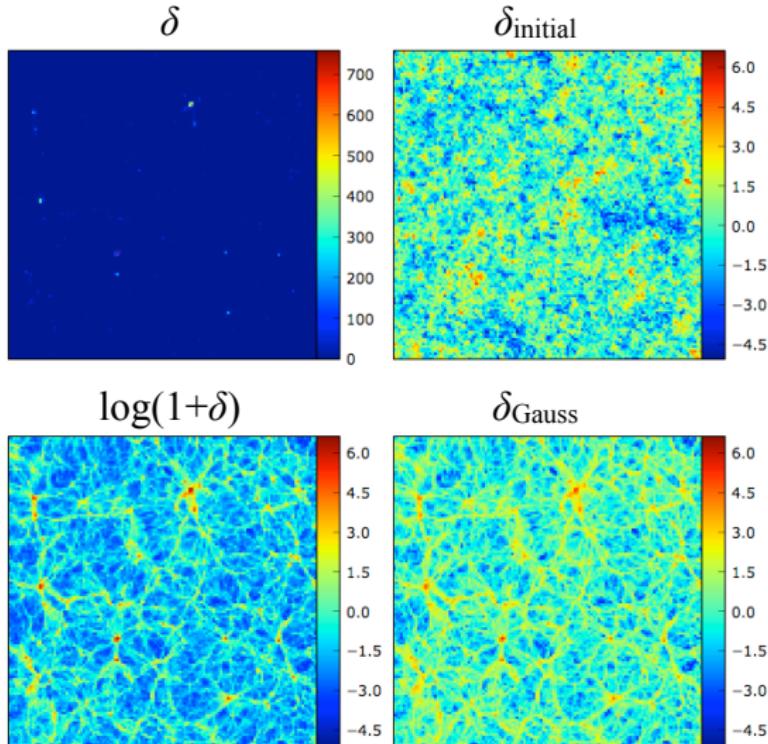
- The standard power spectrum extract a small fraction of the available information from LSS surveys, (unlike CMB)
- Non-Gaussianity, linear modes are used  $k \simeq 0.2$
- Good news: many ( $\propto k^3$ ) high  $k$  modes are available
- Bad news: plateau in the power spectrum due to SS and IS mode coupling
- Standard idea: use higher order statistics (weakly non-linear)
- Worst news: information content of higher order statistics vanishing in the non-linear regime

# Summary of plateaux



# Logarithmic mapping

Neyrinck, Szapudi, & Szalay 2009

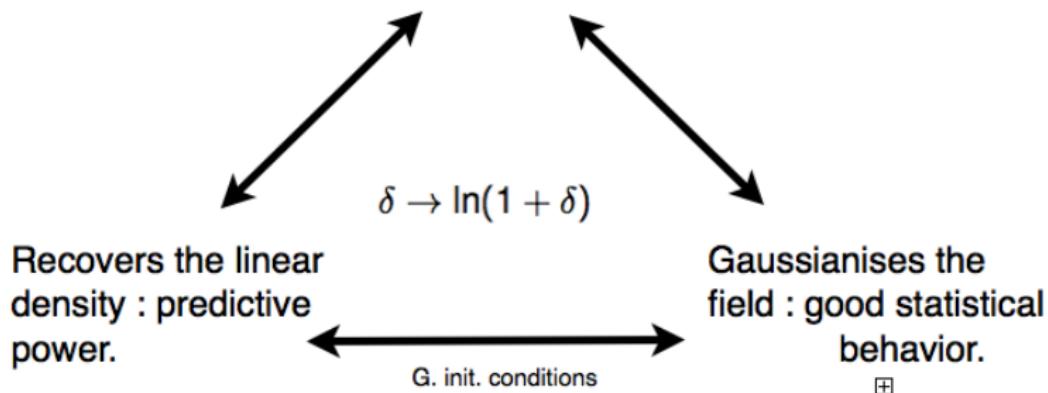


# Sufficient Statistics: All Information on a Parameter

Carron & Szapudi (2013 MNRAS 434, 2961; 2014, MNRAS 439, L11)

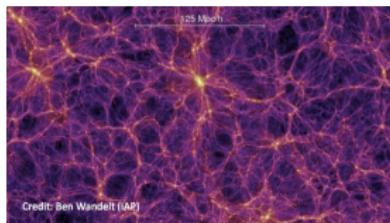
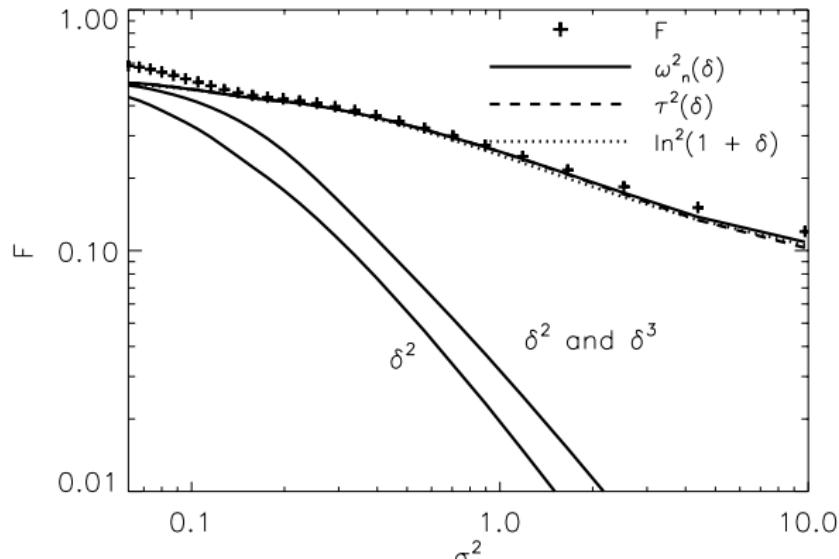
$$\partial_\alpha \ln p(\delta) \simeq \tau^2(\delta) \simeq \left( \frac{(1 + \delta)^{(n+1)/3} - 1}{(n + 1)/3} \right)^2 \simeq \ln(1 + \delta)^2$$

Optimal statistic:  
extracts all information.



Diagonal covariances matrix.

# Info. in the Millenium simulation density field



Analogus to lognormal fields.

## Discreteness effects

So we have now a good understanding at the level of the  $\delta$  field. How to apply these transformation to the observed discrete galaxy field ?

- Now two layers of non-Gaussian statistics : underlying matter field, and discrete sampling effects.

$$P(N|\theta) = \int_0^\infty d\rho \underbrace{p(\rho|\theta)P(N|\rho)}_{\propto P(\rho|N)}.$$

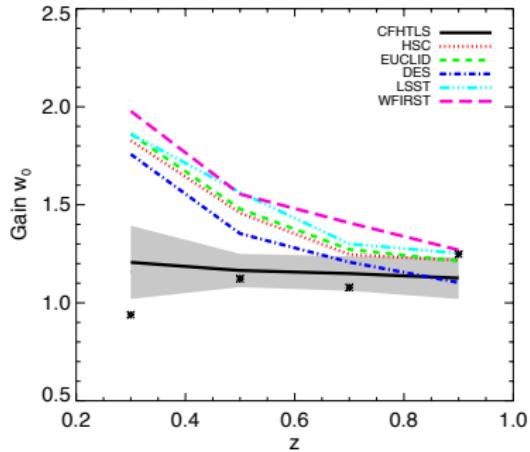
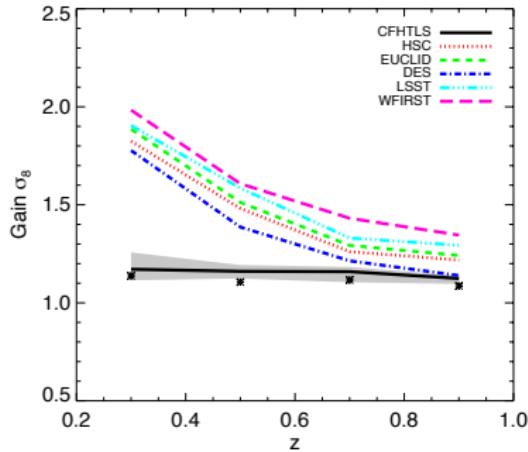
- Saddle-point (Laplace) approximation ( $\rho^*$  maximum of posterior for  $\rho$ ) :

$$\partial_\alpha \ln P(N|\theta) = \underbrace{\partial_\alpha \ln p(\rho^*|\theta)}_{\text{original suff. observable}} - \underbrace{\frac{1}{2} \partial_\alpha \ln'' P(\rho^*|N, \theta)}_{\text{sensitivity of curv. of posterior}} .$$

- For Poisson sampling it all boils down to extract the mean and variance of  $A^* = \ln \rho^*$

# Forecasting dark energy EoS:

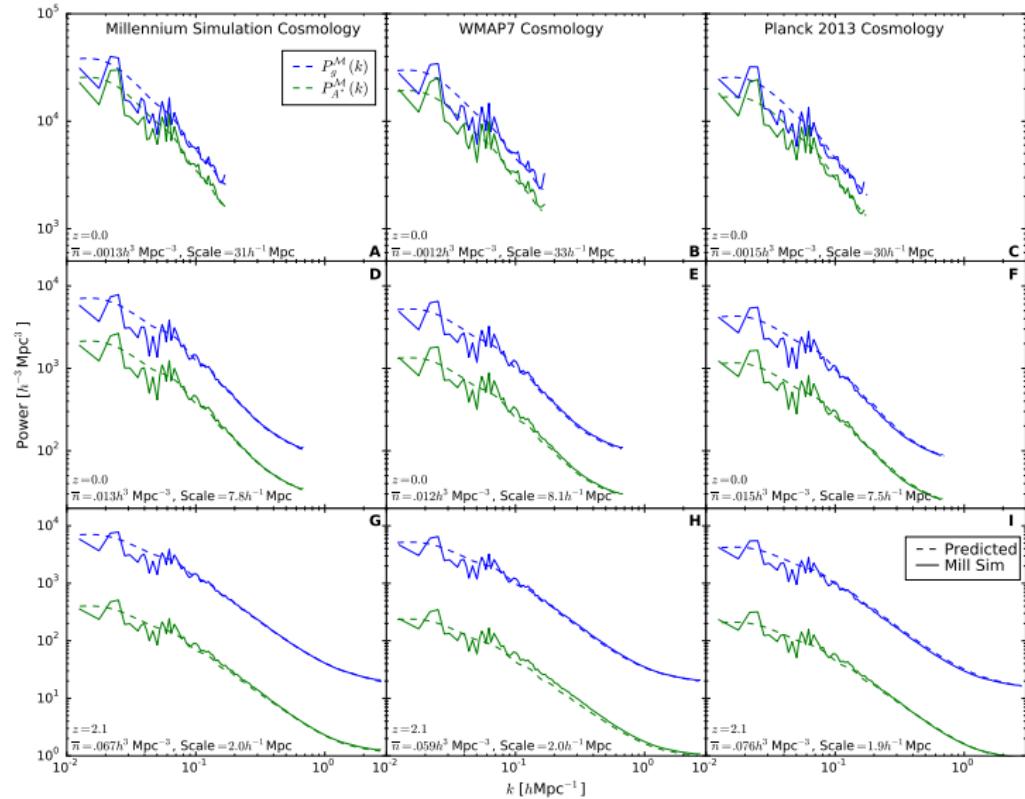
Lognormal is a great approximation in 2D



Survey	Max. Gain
Euclid	1.86
WFIRST-AFTA	1.98
HSC	1.83
LSST	1.86
DES	1.76

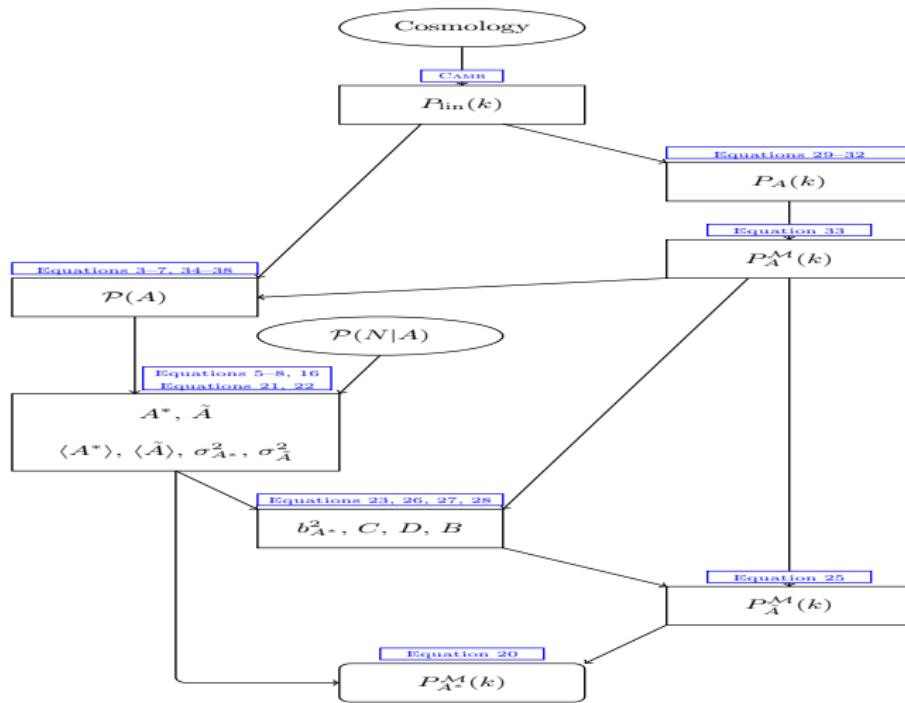
# Prediction of the $A^*$ power spectrum

Repp & Szapudi 2018b, 2019



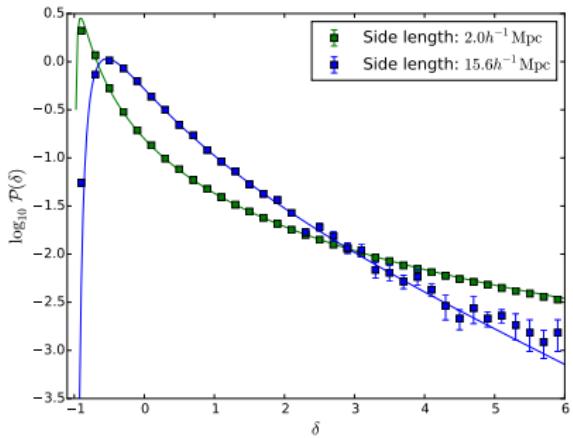
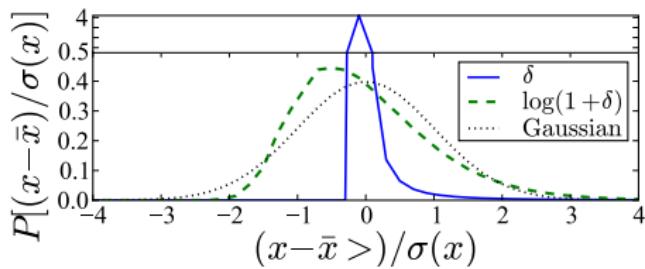
## Prediction of the $A^*$ power spectrum

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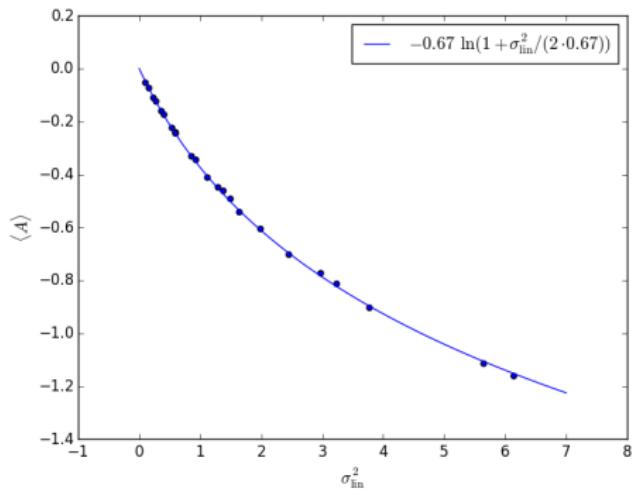
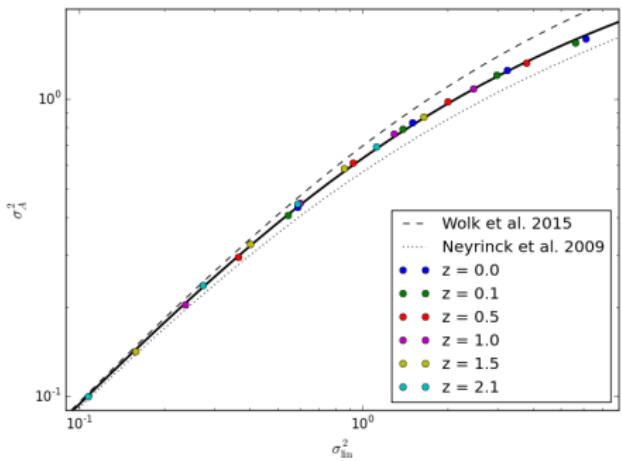


# Non-lognormality in 3D

GEV distribution: Repp & Szapudi 2018a



# Fitting Formulae



# Numerical model for $P_A$

Repp & Szapudi 2017, motivated by Szapudi and Kaiser 2003

- The shape of the  $A = \log(1 + \delta)$  power spectrum is approximately the same as the linear  $\delta$

$$P_A(k) = N C(k) \frac{\sigma_A^2}{\sigma_{\text{lin}}^2} P_{\text{lin}}(k),$$

where

$$C(k) = \begin{cases} 1 & \text{if } k < 0.15 h \text{ Mpc}^{-1} \\ (k/0.15)^\alpha & \text{if } k \geq 0.15 h \text{ Mpc}^{-1} \end{cases},$$

and

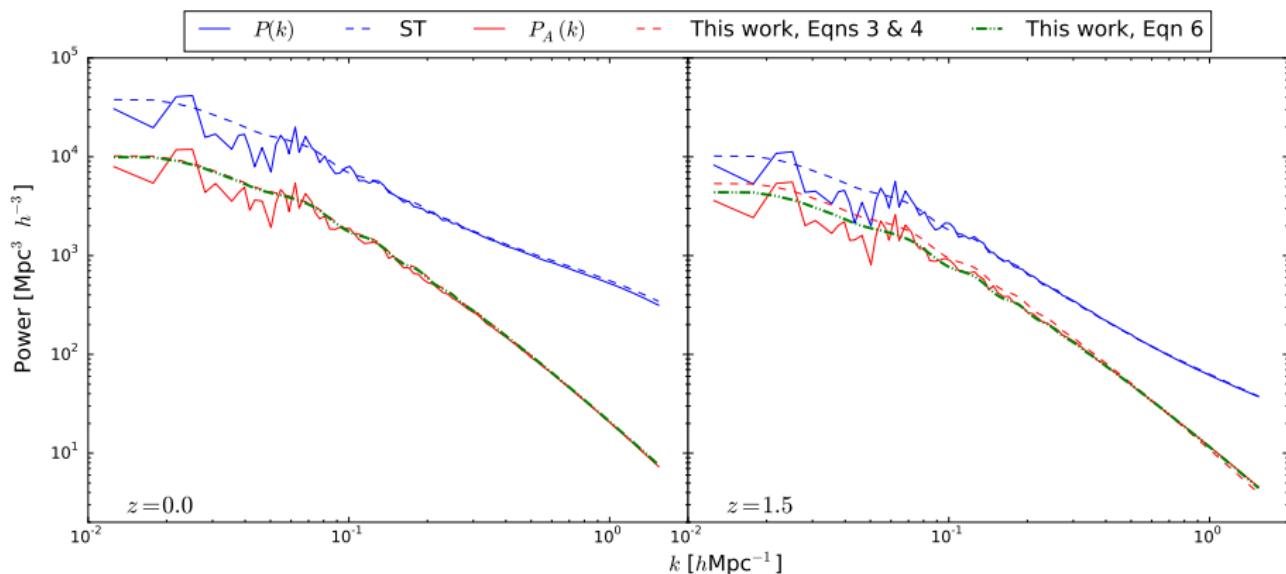
$$N = \frac{\int dk k^2 P_{\text{lin}}(k)}{\int dk k^2 C(k) P_{\text{lin}}(k)}, \quad \alpha(z) \simeq [0.02 - 0.14]$$

- where  $\sigma_{\text{lin}}^2$  is calculated with CAMB and

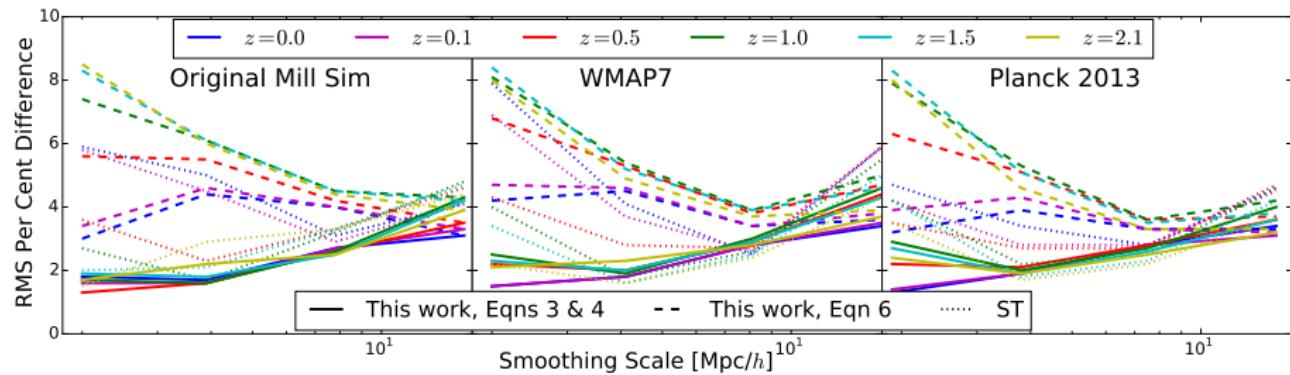
$$\sigma_A^2 = 0.73 \ln \left( 1 + \frac{\sigma_{\text{lin}}^2}{0.73} \right).$$

- checked for several cosmologies and redshifts between  $0 \leq z \leq 2.1$

# Precision Prediction for the Log Power Spectrum

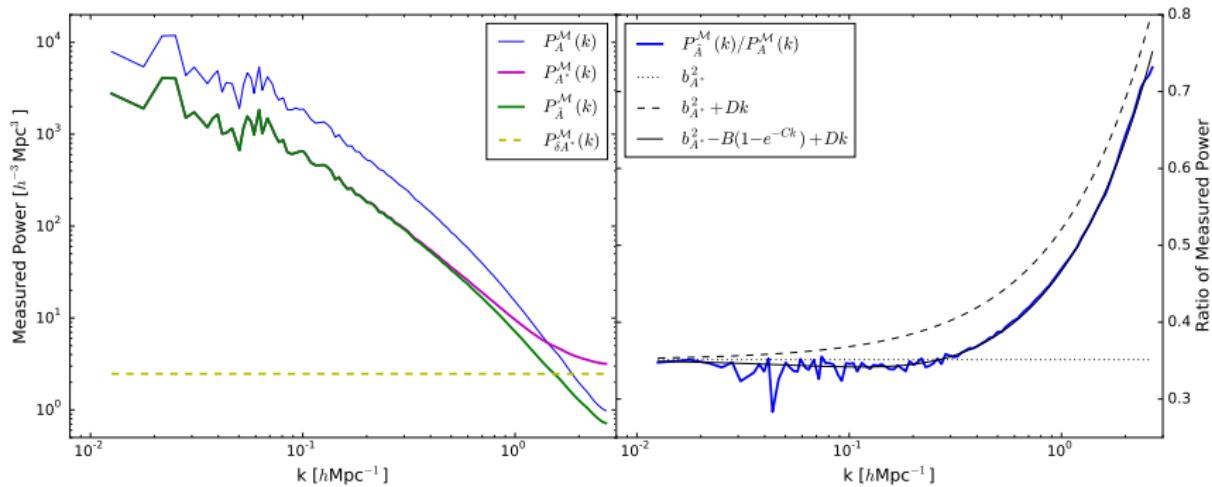


# Prediction errors



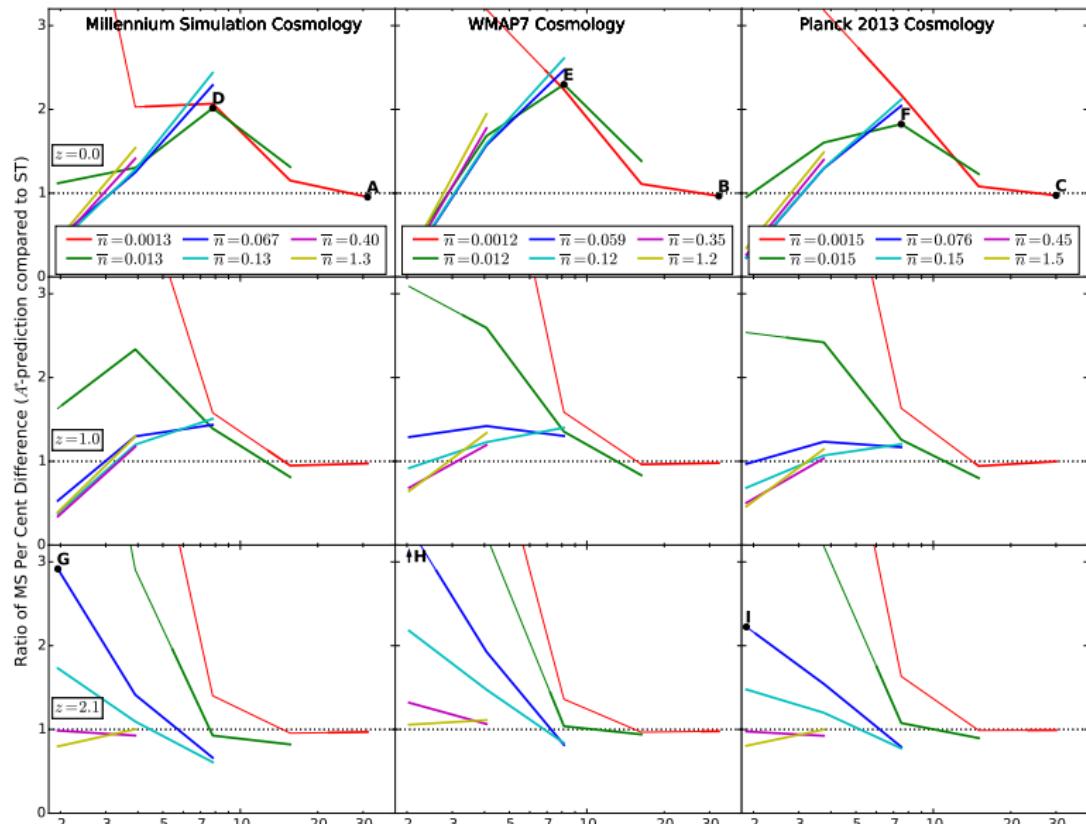
# From $\tilde{A}$ to $A^*$

Repp & Szapudi 2019



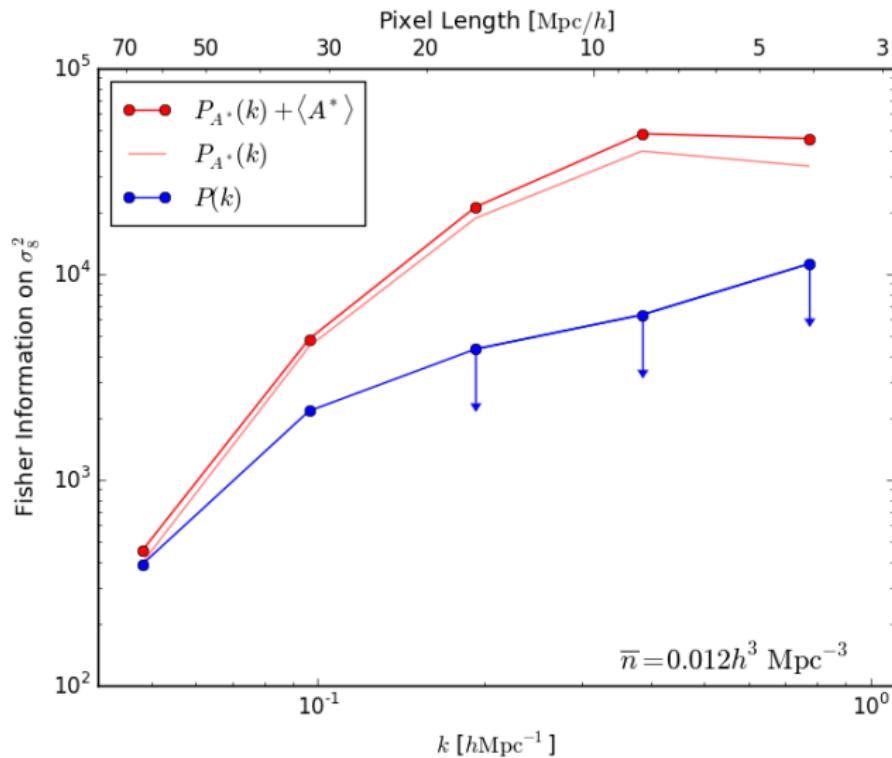
# Accuracy

Repp & Szapudi 2019



# Forecast for $A^*$ in 3D (preliminary)

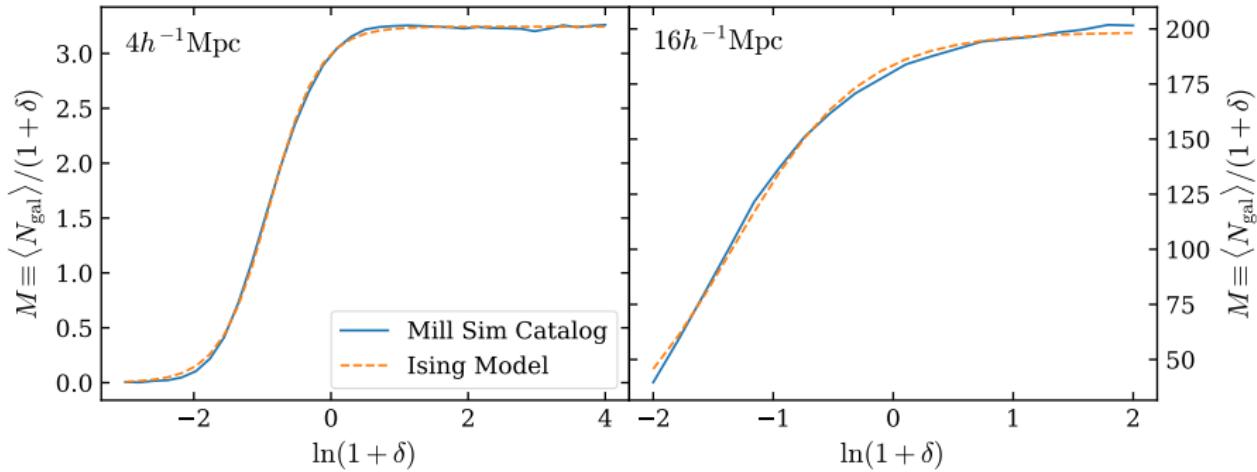
Realistic density but real space no bias



# Bias in Simulations

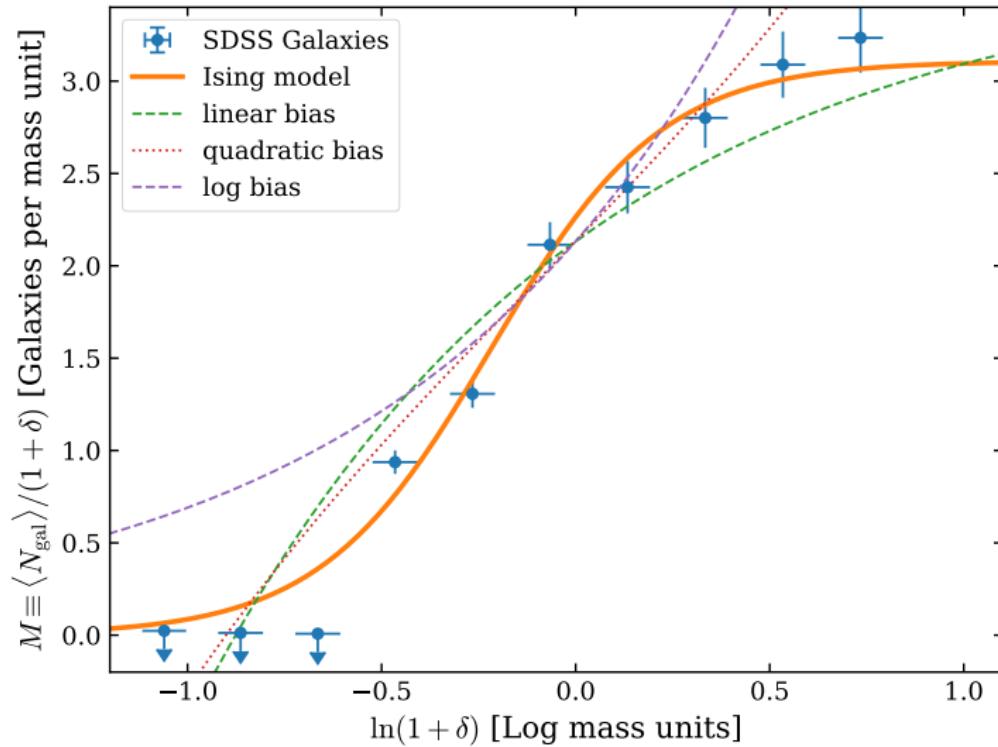
Repp & Szapudi 2019

$$M \equiv \langle N_{\text{gal}} \rangle_A \cdot (1 + \delta)^{-1} = \frac{b\bar{N}}{1 + \exp\left(\frac{A_t - A}{T}\right)}$$

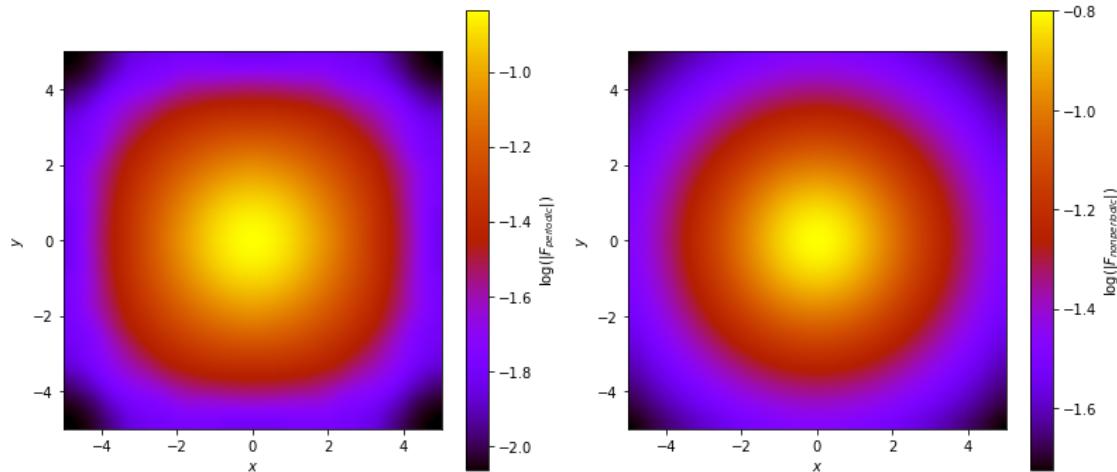


# Bias in SDSS

Repp & Szapudi 2019



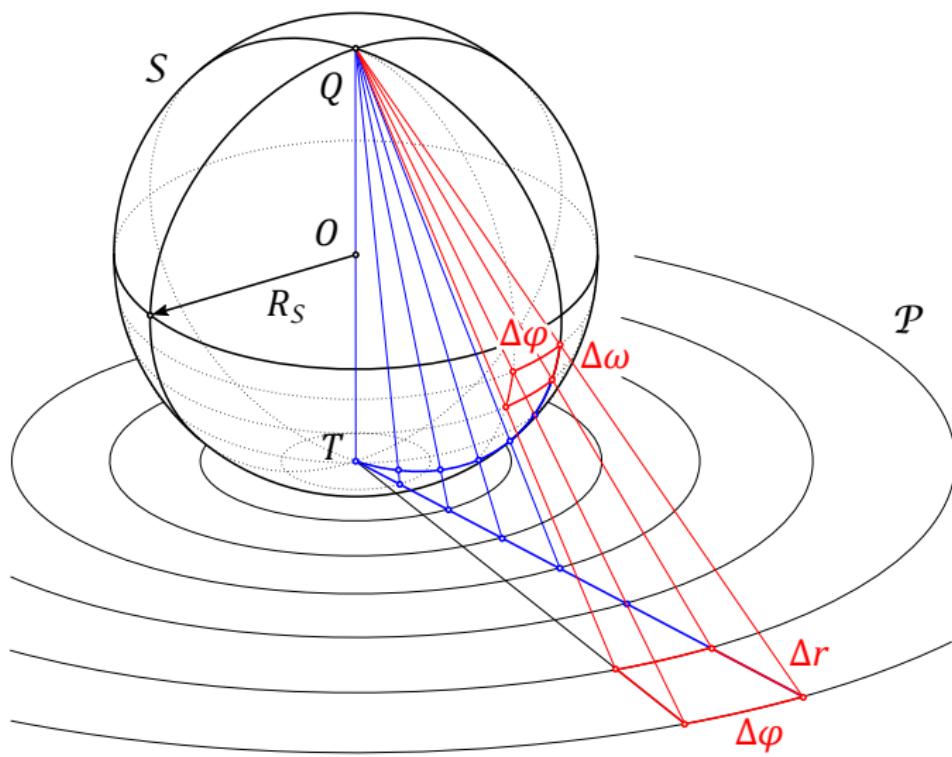
# Problems with cosmological simulations



- Forces are wrong. (High Precision Cosmology?)
- Spherical symmetry broken
- $T^3$  Topology contrary to observations
- Number of modes  $\simeq k^3$ , too many high  $k$
- Does not match observational geometry

# Idea: compactify the infinite volume

Illustrated in 2D

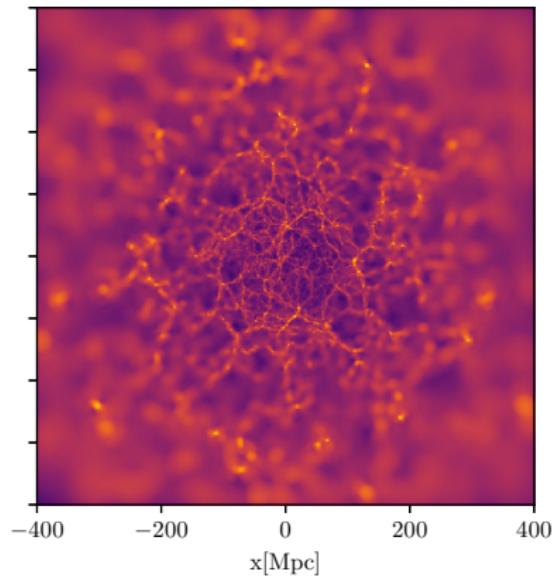
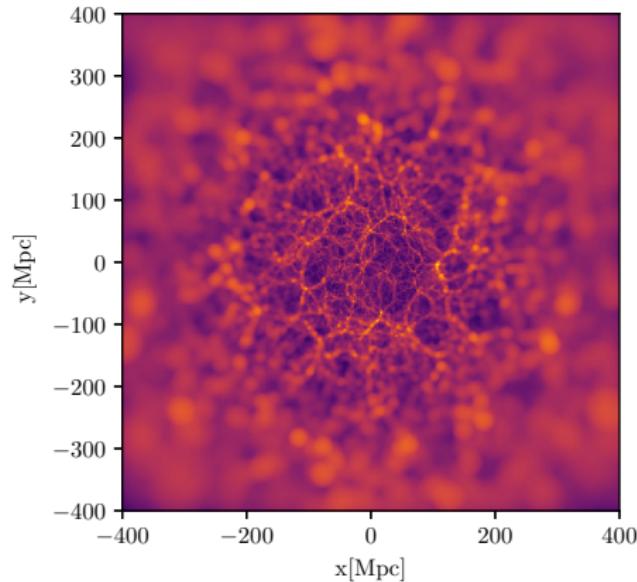


# Millennium resimulations

Racz et al 2019

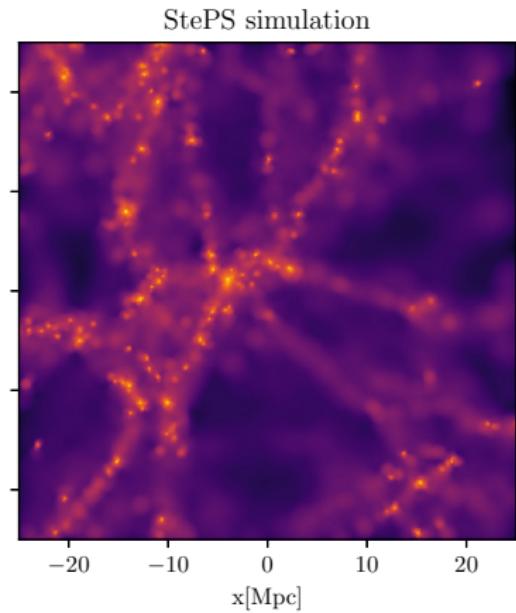
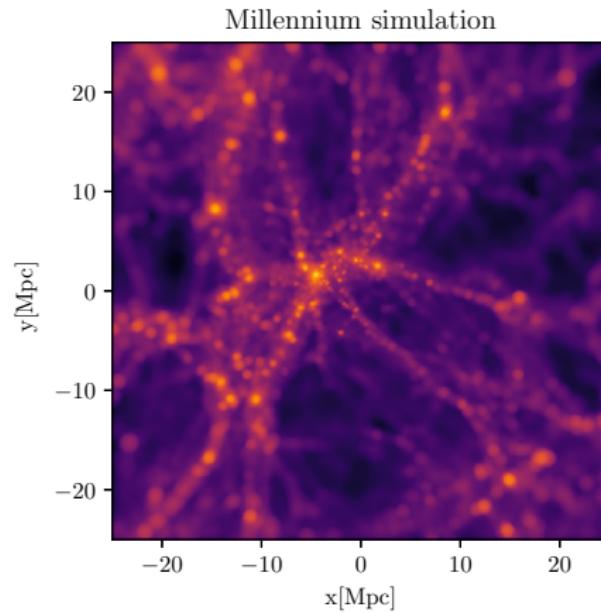
- Millennium Cosmology using the original scripts with StEPS
- $H_0 = 73 \text{ km/s/Mpc}$ ,  $\Omega_m = 0.25$ ,  $\Omega_\Lambda = 0.75$ ,  $\sigma_8 = 0.9$
- 27 volumes to simulate periodic B.C.
- $z_{in} = 127$
- Original:  $512 \text{ CPU} \times 683 \text{ hours } 0.321 \text{ Gpc}^3$
- StEPS:  $12 \text{ GPU} \times 106 \text{ hours } 1.35 \text{ Gpc}^3$
- $N \log N$  vs  $N^2$
- multiresolution

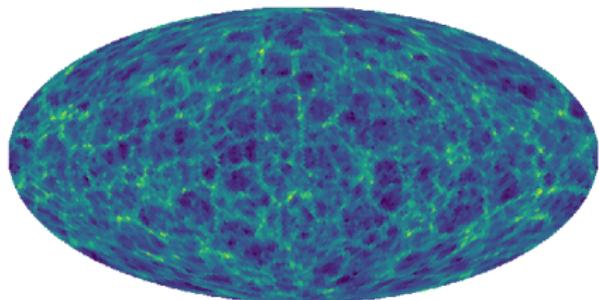
# Results: Millennium



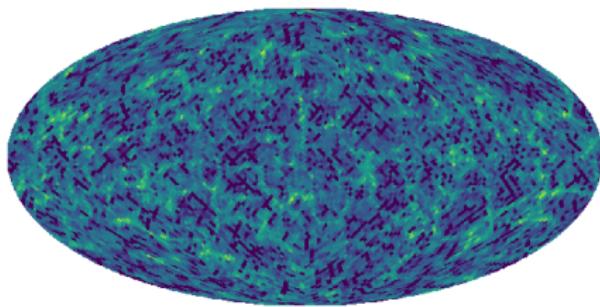
# Results

Racz et al 2019, Millennium resimulation

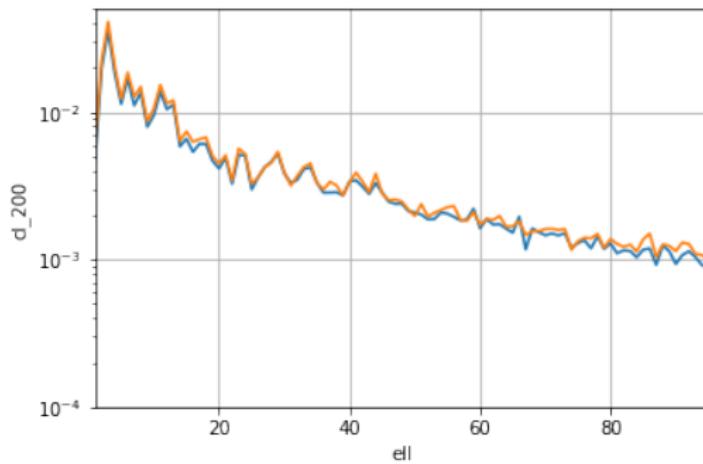




-1.2       $\log(\delta + 1)$       1.8



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## Summary

- Sufficient statistics are well approximated with the log transform
- They are well understood in Fisher information theory, and can be predicted and estimated from data
- The log and  $A^*$  power spectra are now predicted to percent level accuracy in real space
- The information gain for Euclid like surveys  $\gtrsim 2\times$ .
- Accurate bias fit inspired by the Ising model
- Compactified cosmological simulation have unprecedented dynamic range for given resources
- Status: tested on 48 GPUs and recovered Millennium as a test
- Fast prediction for parameter estimation
- Super-survey modes/covariance matrices/high dynamic range (e.g., local group, galaxy simulations in cosmological background)
- Next: redshift distortions

# RSD (preiminary)

