# Towards the analysis of the redshift-space bispectrum

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## **3D galaxy map**



#### **Sloan Digital Sky Survey**

52

цß

 $\mathcal{U}_{O}$ 

65

20h

## **SDSS DR12 BOSS (2016)**

#### **Dard energy:** $w = -1.01 \pm 0.06$

12h

цß

#### Growth rate(z=0.5): $0.452 \pm 0.058$

### Total neutrino mass: $< 0.16 \,\mathrm{eV}/c^2$

Consistent with LCDM. Strong upper limit on total neutrino mass.

 $\mathcal{U}_{O}$ 

70h

## **Next Generation Galaxy Surveys**



## SuMIRe HSC/PFS (2015-25)





#### DESI (4m, LBL, 2020-) LSST (6.5m, SLAC, 2022-)



Euclid (ESA, 2022-)



WFIRST (NASA,2025-)

#### Limitation of two-point statistics



#### **Towards full information extraction**



 $q_0$ 

**Sloan Digital Sky Survey** 

## History

#### Angular catalogues: Peebles & Groth (1975)

Groth & Peebles (1977); Fry & Slender (1982)

#### **Spectroscopic surveys**

#### in configuration-space

Kayo et al. (2004); Jing & Boerner (2004); Wang et al. (2004); Gaztanaga et al. (2005); Nichol et al. (2006); Kulkarni et al. (2007); Gaztanaga et al. (2007); McBride et al. (2011a, b); Marin (2011); Marin et al. (2013); Guo et al. (2013); Slepian et al. (2017a,b);

#### **Spectroscopic surveys in Fourier-space**

Scoccimarro et al. 2001; Feldman et al. 2001; Verde et al. 2002; Gil-Marin et al. 2015a,b;

#### Gil-Marin et al. 2017 (SDSS BOSS analysis)

Pearson & Samushia 2017;

#### Joint analysis of P + B will become the standard method for future galaxy surveys.

## Is there anything else we should do?

Anisotropic bispectrum analysis has not been done yet.

## **Bispectrum Project**

#### [1] NS, Saito, Beutler and Seo 2018

- A new decomposition formalism
- Detection of the quadrupole bispectrum (14 $\sigma$ )

#### [2] In progress

- Modeling the bispectrum
- Modeling the bispectrum covariance
- Fisher analysis

#### [3] Future works

- Analysis using BOSS data
- **Final goal** 
  - Application to future galaxy survey,
     PFS, DESI and Euclid.

## **Decomposition formalism** $B(k_1, k_2, \hat{n})$ Wavevectors Line-of-sight k1 is the z-axis LOS is the z-axis $B = B^m_{\ell_1 \ell_2} Y^m_{\ell_1}(\hat{k}_1) Y^{m*}_{\ell_2}(\hat{k}_2)$ $B = B_{LM} Y_{LM}(\hat{n})$ $\begin{array}{c} (1) \\ \vec{k_2} \end{array} \stackrel{\widetilde{\mathbf{h}}}{\stackrel{\phantom{}}{\mathbf{h}}} \vec{k_1} \\ \overrightarrow{\mathbf{h}} \vec{n} \end{array}$ $\hat{y}$ $arphi_{12}$

Scoccimarro et al. (1999)

Slepian et al. (2017)

#### New decomposition formalism not depending on coordinate systems

1)Expand the bispectrum in three spherical harmonics

$$\begin{split} B^{m_1m_2M}_{\ell_1\ell_2L}(k_1,k_2) &= N_{\ell_1\ell_2L} \int \frac{d^2\hat{k}_1}{4\pi} \int \frac{d^2\hat{k}_2}{4\pi} \int \frac{d^2\hat{n}}{4\pi} \\ &\times y^{m_1*}_{\ell_1}(\hat{k}_1) y^{m_2*}_{\ell_2}(\hat{k}_2) y^{M*}_L(\hat{n}) B(k_1,k_2,\hat{n}), \end{split}$$

2) Sum up over all m-modeswith wigner 3j symbol.

$$B_{\ell_1\ell_2L}(k_1,k_2) \propto \sum_{m_1m_2m_3} \left( \begin{smallmatrix} \ell_1 & \ell_2 & \ell_3 \\ m_1 & m_2 & m_3 \end{smallmatrix} \right) B_{\ell_1\ell_2L}^{m_1m_2M}(k_1,k_2)$$

3) Restrict the allowed multipoles to I1+I2+L=even

#### **New Bispectrum Multipoles** Three multipole indexes Two wavenumber



#### New Bispectrum Multipoles Three multipole indexes Two wavenumber



## Advantages (1)

#### For example,

$$B_{\ell_1 \ell_2 L}(k_1, k_2) \propto \int \frac{d \cos \theta}{2} \sum_M \begin{pmatrix} \ell_1 & \ell_2 & L \\ 0 & -M & M \end{pmatrix}$$
  
Ours 
$$\times \mathcal{L}_{\ell_2}^{-M}(\cos \theta) B_{LM}(k_1, k_2, \theta)$$

#### Scoccimarro et al. (1999)

## Independent of the choice of the coordinate axis.



#### Following Scoccimarro 2015:

$$\begin{split} \widehat{B}_{\ell_{1}\ell_{2}L}(k_{1},k_{2}) &= H_{\ell_{1}\ell_{2}L} \sum_{m_{1}m_{2}M} \begin{pmatrix} \ell_{1} & \ell_{2} & L \\ m_{1} & m_{2} & M \end{pmatrix} \\ & \times \frac{N_{\ell_{1}\ell_{2}L}}{I} \int d^{3}x \, F_{\ell_{1}}^{m_{1}}(\boldsymbol{x};k_{1}) \, F_{\ell_{2}}^{m_{2}}(\boldsymbol{x};k_{2}) \, G_{L}^{M}(\boldsymbol{x}), \end{split}$$

$$\begin{split} F_{\ell}^{m}(\boldsymbol{x};k) &= \int \frac{d^{2}\hat{k}}{4\pi} \, \mathrm{e}^{\mathrm{i}\boldsymbol{k}\cdot\boldsymbol{x}} \, y_{\ell}^{m*}(\hat{k}) \, \frac{\delta n|_{\mathrm{FFT}}(\boldsymbol{k})}{W_{\mathrm{mass}}(\boldsymbol{k})} \\ G_{L}^{M}(\boldsymbol{x}) &= \int \frac{d^{3}k}{(2\pi)^{3}} \, \mathrm{e}^{\mathrm{i}\boldsymbol{k}\cdot\boldsymbol{x}} \, \frac{\delta n_{L}^{M}|_{\mathrm{FFT}}(\boldsymbol{k})}{W_{\mathrm{mass}}(\boldsymbol{k})}. \end{split}$$

## **Requires only FFT processes.**

#### **Advantages (3)** Double Hankel transform



$$\zeta_{000}^{\text{masked}}(r_{1}, r_{2}) = Q_{000}(r_{1}, r_{2})\zeta_{000}^{\text{theory}}(r_{1}, r_{2}) + Q_{110}(r_{1}, r_{2})\zeta_{110}^{\text{theory}}(r_{1}, r_{2}) + \frac{Q_{110}(r_{1}, r_{2})\zeta_{110}^{\text{theory}}(r_{1}, r_{2})}{\text{Window function}}$$

## Gives a simple expression to correct for survey geometry

#### **Three-point window function multipoles**



## Measurements of Monopole (L=0) (k1 = k2)



## Quadrupole (L=2) and Hexadecapole (L=4)



#### Strategy

#### Focus only on k1=k2 elements of the bispectrum multipoles.

Decreases the number of bins

- Smaller number of mocks
- Fast measurements
  - Fast analysis

2) Full bispectrum analysis.

## **B000**

#### Full(k1, k2)



k1 = k2



## Cumulative S/N: CMASS North



How do anisotropic bispectrum measurements improve constraints on cosmological parameters? (Preliminary results)

## Fisher analysis

## Correct error estimates [non-Gaussian covariance]

## **Correct theoretical model**

## Importance of Non-Gaussian errors



#### **Non-Gaussian covariance**

#### P-P covariance -> tri-spectrum P-B covariance -> 5-point spectrum B-B covariance -> 6-point spectrum

## RSD + linear bias (b1) + shot noise

#### **Shot-noise subtraction effect**

$$Cov (|\delta(\mathbf{k})|^2, |\delta(\mathbf{k})|^2)$$
$$Cov (P(\mathbf{k}), P(\mathbf{k}))$$
$$P = |\delta(\mathbf{k})|^2 - \frac{1}{\bar{n}}$$
$$N_{obs} = N_{true} (1 + \delta_b)$$

## Before shot-noise subtraction Cov [P, P]

$$= \frac{1}{V} \Big\{ T(\boldsymbol{k}, -\boldsymbol{k}, \boldsymbol{k'}, -\boldsymbol{k'}) \Big\}$$

+  $(1/\bar{n})[B(k, -k, 0) + B(k', -k', 0)]$ 

+ B(k, k', -k - k') + B(k, -k, -k + k') + B(-k, k', k - k') + B(-k, -k', k + k')]+  $(1/\bar{n}^2)[2P(k) + 2P(k') + P(0) + P(k + k') + P(k - k')]$ 

+  $(1/\bar{n}^3)$ 

## After shot-noise subtraction Cov [P, P]

$$= \frac{1}{V} \Big\{ T(\boldsymbol{k}, -\boldsymbol{k}, \boldsymbol{k}', -\boldsymbol{k}') \Big\}$$

+  $(1/\bar{n})[B(k, \pi, 0)] + B(k', \pi', 0)$ 

+ B(k, k', -k - k') + B(k, -k, -k + k') + B(-k, k', k - k') + B(-k, -k', k + k')]

- +  $(1/\bar{n}^2)[2P(\mathbf{x}) + 2P(\mathbf{x}) + P(\mathbf{k} + \mathbf{k'}) + P(\mathbf{k} \mathbf{k'})]$ +  $(1/\bar{n}^2)[2P(\mathbf{x}) + 2P(\mathbf{x}) + P(\mathbf{k} \mathbf{k'})]$

#### **Shot-noise subtraction effect**

$$\operatorname{Cov}(B, B)$$

$$\operatorname{Cov}(\delta^{3}, \delta^{3})$$

$$B = \delta^3 - \frac{1}{\bar{n}}P - \frac{1}{\bar{n}^2}$$

## Cov [PO, PO]



## **Cov** [PO, B000]



## **Cov [B000, B000]**



**S**/**N** Monopole(B000) Quadrupole(B202)



The same signal is used for both the mock and PT.

### **Bispectrum model**

### We need a model that explains the anisotropic galaxy power spectrum and bispectrum simultaneously.

## Non-linear clustering Non-linear RSDs Non-linear bias

### **Bispectrum model**

## We need a model that explains the anisotropic ( ) wer spectrum and bispectrum nultaneously.

Non-Imar clustering
 Non-linear RSDs
 Non-linear bias

## Bispectrum model describing the BAO damping In the power spectrum case (Eisenstein+2007):

$$P(\boldsymbol{k}, \hat{n}) = \left[ Z^{[1]}(\boldsymbol{k}, \hat{n}) \right]^2 \left[ \mathcal{D}^2(\boldsymbol{k}, \hat{n}) P_{\text{BAO}}(\boldsymbol{k}) + P_{\text{nw}}(\boldsymbol{k}) \right]$$
  
BAO Power spec.  
without BAO

$$\begin{aligned} \mathcal{D}(\boldsymbol{k}, \hat{n}) &= \exp\left(-\frac{k^2 \mu^2 \Sigma_{\parallel}^2 + k^2 \left(1 - \mu^2\right) \Sigma_{\perp}^2}{4}\right) \\ P_{\text{BAO}}(k) &= P_{\text{lin}}(k) - P_{\text{nw}}(k) \end{aligned}$$



## **IR cancellation**

## Assuming that the IR flow is NOT

#### correlated with the density field:

$$egin{aligned} &\langle \delta(oldsymbol{x}-\overline{oldsymbol{\Psi}}) &\langle \delta(oldsymbol{x}-\overline{oldsymbol{\Psi}}) 
angle &= \langle \delta(oldsymbol{x}) \delta(oldsymbol{y}) 
angle \ &= \xi(oldsymbol{x}-oldsymbol{y}) \end{aligned}$$

#### The effect from the IR flow completely cancels out because of translational symmetry (Galilean invariance)

## **IR cancellation**

## Assuming that the IR flow is NOT

#### correlated with the density field:

$$\langle e^{-i\boldsymbol{k}_{1}\cdot\overline{\boldsymbol{\Psi}}}\widetilde{\delta}(\boldsymbol{k}_{1})e^{-i\boldsymbol{k}_{2}\cdot\overline{\boldsymbol{\Psi}}}\widetilde{\delta}(\boldsymbol{k}_{2})\rangle$$

$$= \langle e^{-i\boldsymbol{k}_{1}\cdot\overline{\boldsymbol{\Psi}}}e^{-i\boldsymbol{k}_{2}\cdot\overline{\boldsymbol{\Psi}}}\rangle\langle\widetilde{\delta}(\boldsymbol{k}_{1})\widetilde{\delta}(\boldsymbol{k}_{2})\rangle$$

$$= (2\pi)^{3}\delta_{\mathrm{D}}\left(\boldsymbol{k}_{1}+\boldsymbol{k}_{2}\right)P(\boldsymbol{k}_{1})$$

#### The effect from the IR flow completely cancels out because of translational symmetry (Galilean invariance)

## **IR (high-k limit) cancellation**

$$P_{22} = \int \frac{d^3 p_1}{(2\pi)^3} \int \frac{d^3 p_2}{(2\pi)^3} \times (2\pi)^3 \delta_{\mathrm{D}} (\mathbf{k} - \mathbf{p}_1 - \mathbf{p}_2) F^2(\mathbf{p}_1, \mathbf{p}_2) P_{\mathrm{lin}}(p_1) P_{\mathrm{lin}}(p_2)$$

$$p_1 \gg p_2 \quad \text{Shift}$$

$$\rightarrow k^2 \sigma^2 P_{\mathrm{lin}}(k)$$

$$P_{13} + P_{22} \to 0$$

**IR (high-k limit) cancellation Gamma-expansion (all orders)**  $P(\boldsymbol{k}) = G^{2}(\boldsymbol{k}) P_{\text{lin}}(\boldsymbol{k}) + P_{\text{MC}}(\boldsymbol{k})$  $\mathcal{D}^2(\mathbf{k})[Z^{[1]}(\mathbf{k})]^2 P_{\mathrm{lin}}(k)$  $\left[1 - \mathcal{D}^2(\boldsymbol{k})\right] \left[Z^{[1]}(\boldsymbol{k})\right]^2 P_{\text{lin}}(\boldsymbol{k})$  $\rightarrow [Z^{[1]}(\boldsymbol{k})]^2 P_{\text{lin}}(k)$ 

The IR flow never affects the power spectrum??

To extract physical effects from the IR flow, the breaking of the IR cancellation should be considered.

## Breaking of the IR (high-k limit) cancellation

$$P_{22} = \int \frac{d^3 p_1}{(2\pi)^3} \int \frac{d^3 p_2}{(2\pi)^3}$$

$$\times (2\pi)^3 \delta_{\mathrm{D}} (\mathbf{k} - \mathbf{p}_1 - \mathbf{p}_2) F^2(\mathbf{p}_1, \mathbf{p}_2) P_{\mathrm{lin}}(p_1) P_{\mathrm{lin}}(p_2)$$

$$p_1 \gg p_2 \qquad \qquad \mathbf{Too strong}$$

$$\rightarrow k^2 \sigma^2 P_{\mathrm{lin}}(k)$$

## Breaking of the IR (high-k limit) cancellation



Eisenstein's template  

$$P(\mathbf{k}) = G^{2}(\mathbf{k}) P_{\text{lin}}(\mathbf{k}) + P_{\text{MC}}(\mathbf{k})$$

$$\mathcal{D}^{2}(\mathbf{k})[Z^{[1]}(\mathbf{k})]^{2}P_{\text{lin}}(\mathbf{k})$$

$$[1 - \mathcal{D}^{2}(\mathbf{k})] [Z^{[1]}(\mathbf{k})]^{2}P_{n^{W}}(\mathbf{k})$$

$$P(\boldsymbol{k},\hat{n}) = \left[Z^{[1]}(\boldsymbol{k},\hat{n})\right]^2 \left[\mathcal{D}^2(\boldsymbol{k},\hat{n}) P_{\text{BAO}}(\boldsymbol{k}) + P_{\text{nw}}(\boldsymbol{k})\right]$$



 $\mathcal{D}(oldsymbol{k}_1)\mathcal{D}(oldsymbol{k}_2)\mathcal{D}(oldsymbol{k}_{12})B_{ ext{tree}}(oldsymbol{k}_1,oldsymbol{k}_2) = \left(\mathcal{D}^2(oldsymbol{k}_1)-\mathcal{D}(oldsymbol{k}_1)\mathcal{D}(oldsymbol{k}_2)\mathcal{D}(oldsymbol{k}_{12})
ight)B_{ ext{tree}}(oldsymbol{k}_1,oldsymbol{k}_2) = \left(\mathcal{D}^2(oldsymbol{k}_1)-\mathcal{D}(oldsymbol{k}_1)\mathcal{D}(oldsymbol{k}_2)\mathcal{D}(oldsymbol{k}_{12})
ight)B_{ ext{tree}}(oldsymbol{k}_1,oldsymbol{k}_2)$ 



 $egin{aligned} & \left(\mathcal{D}(m{k}_1)\mathcal{D}(m{k}_2)\mathcal{D}^{-1}(m{k}_{12}) - \mathcal{D}^2(m{k}_1) - \mathcal{D}^2(m{k}_2) 
ight. \ & +\mathcal{D}(m{k}_1)\mathcal{D}(m{k}_2)\mathcal{D}(m{k}_{12})
ight) B_{ ext{tree}}(m{k}_1,m{k}_2) \end{aligned}$ 



$$(1 - \mathcal{D}(\boldsymbol{k}_1)\mathcal{D}(\boldsymbol{k}_2)\mathcal{D}^{-1}(\boldsymbol{k}_{12}))B_{\text{tree}}(\boldsymbol{k}_1, \boldsymbol{k}_2)$$





 $\mathcal{D}(oldsymbol{k}_1)\mathcal{D}(oldsymbol{k}_2)\mathcal{D}(oldsymbol{k}_{12})B_{ ext{tree}}(oldsymbol{k}_1,oldsymbol{k}_2) = \left(\mathcal{D}^2(oldsymbol{k}_1)-\mathcal{D}(oldsymbol{k}_1)\mathcal{D}(oldsymbol{k}_2)\mathcal{D}(oldsymbol{k}_{12})
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ight)B_{ ext{tree}}(oldsymbol{k}_1,oldsymbol{k}_2)$ 





$$(1 - \mathcal{D}(\boldsymbol{k}_1)\mathcal{D}(\boldsymbol{k}_2)\mathcal{D}^{-1}(\boldsymbol{k}_{12}))B_{\text{tree}}(\boldsymbol{k}_1, \boldsymbol{k}_2)$$



### Model

$$B(\mathbf{k}_{1}, \mathbf{k}_{2}) = 2 Z^{[1]}(\mathbf{k}_{1}) Z^{[1]}(\mathbf{k}_{2}) Z^{[2]}(\mathbf{k}_{1}, \mathbf{k}_{2}) \times \left\{ \mathcal{D}(\mathbf{k}_{1}) \mathcal{D}(\mathbf{k}_{2}) \mathcal{D}(\mathbf{k}_{12}) P_{\text{BAO}}(k_{1}) P_{\text{BAO}}(k_{2}) + \mathcal{D}^{2}(\mathbf{k}_{1}) P_{\text{BAO}}(k_{1}) P_{\text{nw}}(k_{2}) + \mathcal{D}^{2}(\mathbf{k}_{2}) P_{\text{nw}}(k_{1}) P_{\text{BAO}}(k_{2}) + P_{\text{nw}}(k_{1}) P_{\text{nw}}(k_{2}) \right\} + 2 \text{ cyc.}$$

$$\mathcal{D}(\boldsymbol{k}, \hat{n}) = \exp\left(-\frac{k^2 \mu^2 \Sigma_{\parallel}^2 + k^2 \left(1 - \mu^2\right) \Sigma_{\perp}^2}{4}\right)$$
$$P_{\text{BAO}}(k) = P_{\text{lin}}(k) - P_{\text{nw}}(k)$$

#### **BAO damping**

$$B_{000}(k,k) = B_{BAO,000}(k,k) + B_{nw,000}(k,k) \left[ak + bk^2 + ck^3\right]$$



**Bispectrum model for the fisher analysis** 

- Power spectrum and bispectrum: Tree-level
  - + non-linear BAO damping

  - + local bias (b1, b2) <-- non-linear bias
  - + residual shot-noise terms



## PFS Project (0.6 < z < 2.2)



## **Bispectrum Project**

#### [1] NS, Saito, Beutler and Seo 2018

- A new decomposition formalism
- Detection of the quadrupole bispectrum (14 $\sigma$ )

This talk

#### [2] In progress

- Modeling the bispectrum covariance
- Fisher analysis
- Modeling the anisotropic bispectrum

#### [3] Future works

Analysis using BOSS data

#### **Final goal**

Application to future galaxy survey,
 PFS, DESI and Euclid.